

# Magnetic properties of Mn-doped $\text{Bi}_2\text{Se}_3$ compound: temperature dependence and pressure effects

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## Abstract

Magnetic susceptibility  $\chi$  of  $\text{Bi}_{2-x}\text{Mn}_x\text{Se}_3$  ( $x = 0.01-0.2$ ) was measured in the temperature range 4.2–300 K. For all the samples, a Curie–Weiss behaviour of  $\chi(T)$  was revealed with effective magnetic moments of Mn ions corresponding to the spin value  $S = 5/2$ , which couple antiferromagnetically with the paramagnetic Curie temperature  $\Theta \sim -50$  K. In addition, for the samples of nominal composition  $x = 0.1$  and  $0.2$  the effect of a hydrostatic pressure  $P$  up to 2 kbar on  $\chi$  has been measured at fixed temperatures 78 and 300 K that allowed to estimate the pressure derivative of  $\Theta$  to be  $d\Theta/dP \sim -0.8$  K kbar<sup>-1</sup>. Based on the observed behaviour of  $\Theta$  with varied Mn concentration and pressure, a possible mechanism of interaction between localized Mn moments is discussed.

Keywords: diluted magnetic semiconductors,  $\text{Bi}_2\text{Se}_3$ , pressure effects, paramagnetic curie temperature

(Some figures may appear in colour only in the online journal)

## 1. Introduction

New class of tetradymite semiconductors  $\text{Bi}_2\text{Te}_3$ ,  $\text{Bi}_2\text{Se}_3$  and  $\text{Sb}_2\text{Te}_3$  is of great attraction due to a variety of unconventional transport properties related to an anomalous band structure that supports topologically gapless surface states of the Dirac cone type (see, e.g. [1–3] and references therein). Of particular interest are the effects of magnetic impurities and ferromagnetism on the surface states [4]. The magnetic dopants such as Fe or Mn are expected to open the gap at the Dirac point [4–6] that would provide a doping control of the many topological phenomena and could lead to their unique practical applications in new spintronic or magnetoelectric devices. In this context, it is of great importance a detailed study of the magnetic properties of these compounds to shed more light on the nature of magnetic interaction between dopant magnetic moments and the interplay between their surface and bulk magnetism.

Focusing here on the bulk properties, we would like to mention a number of experimental results on bulk magnetism which has been reported for magnetically doped topological insulators:  $\text{Bi}_{2-x}\text{M}_x\text{Te}_3$  ( $M = \text{Fe}$  [7, 8], Mn [9–12]),  $\text{Bi}_{2-x}\text{M}_x\text{Se}_3$  ( $M = \text{Fe}$  [13, 14], Mn [12, 15, 16], Cr [14]),  $\text{Sb}_{2-x}\text{M}_x\text{Te}_3$  ( $M = \text{Fe}$  [17], Mn [12, 18, 19], Cr [20], V [21, 22]). Another direction of research concerns experimental studies of peculiar magnetic properties of Mn-doped  $\text{Bi}_2\text{Se}_3$  thin films, which were carried out in the recent years (see detailed papers [23–25] and references therein). The obtained results revealed a wide variety of magnetic properties depending on the type of the 3d-dopant as well as the host compound. For example, the  $\text{Bi}_{2-x}\text{Fe}_x\text{Te}_3$  compounds demonstrate a ferromagnetic ordering with low-spin state of Fe ions [7, 8] while in  $\text{Sb}_{2-x}\text{Fe}_x\text{Te}_3$  compounds Fe ions are in a high-spin state and interact antiferromagnetically [17]. Very different behaviour of the magnetic dopant in topological insulators was also confirmed by theoretical calculations for some considered

systems [5, 26, 27]. Despite the substantial amount of available information on the bulk magnetism of these compounds, the inconsistency of some experimental data on the magnetic state of 3d-dopants [13, 14] and character of their interaction [11, 18] should be emphasized. Furthermore, there is no unified viewpoint on the mechanism of interaction between the magnetic impurities, and the proposed theoretical models of the exchange interaction [27–31] need to be verified by experiment.

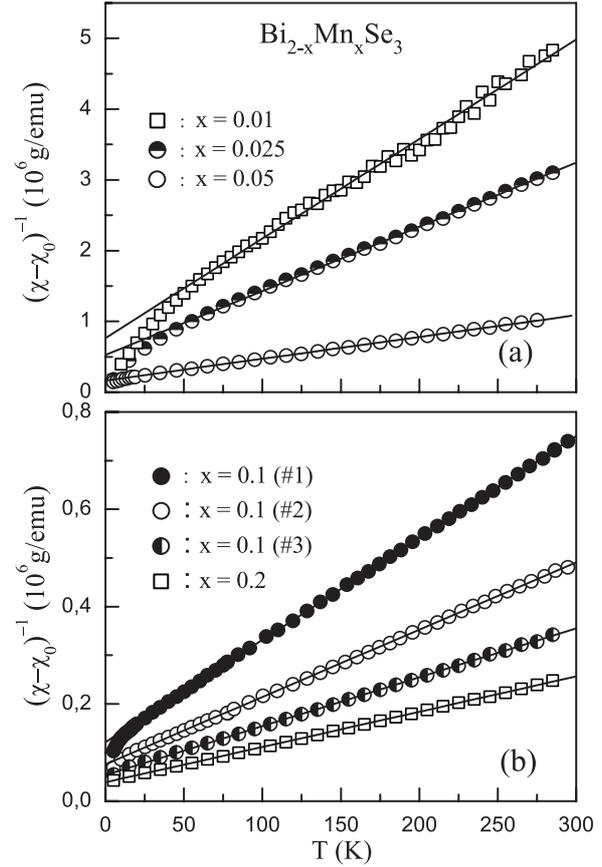
In this paper we have attempted to contribute to the problem of diluted magnetic semiconductors by studying the bulk magnetic properties of  $\text{Bi}_{2-x}\text{Mn}_x\text{Se}_3$  system in a wide concentration range ( $0 \leq x \leq 0.2$ ). The aim of the study was to clarify the magnetic parameters of the system, as well as to reveal specific features in their concentration dependence at the limits of solubility of manganese. In addition, the effects of high pressure on the magnetic susceptibility of the system were measured here for the first time. These measurements were motivated by the huge pressure effects on the transport properties observed in the pure  $\text{Bi}_2\text{Se}_3$  compound [32, 33]. It should be pointed out that along with the widely used doping methods, pressure is a particularly powerful tool for tuning the electronic properties of solids without introducing a disorder inherent to the chemical substitution.

## 2. Experimental

Single crystals of  $\text{Bi}_{2-x}\text{Mn}_x\text{Se}_3$  with nominal compositions of  $x = 0.01, 0.025, 0.05, 0.10$  and  $0.20$  were grown from a melt using the Bridgman method. Corresponding amounts of Bi, Mn, Se of minimum purity 99.99% were mixed and sealed in evacuated silica ampoules. The ampoules were annealed at  $850^\circ\text{C}$  for 16h for homogenization. The melt was then cooled down to  $550^\circ\text{C}$  at a rate of  $4^\circ\text{C h}^{-1}$  and then down to room temperature with a rate of  $100^\circ\text{C h}^{-1}$ . Well-formed silvery crystal rods were obtained, which could be easily cleaved into plates with flat shiny surfaces. Phase purity of the obtained samples was characterized at room temperature by powder x-ray diffraction (XRD) using a D8 Advance Bruker AXS diffractometer with  $\text{Cu K}\alpha$  radiation. XRD experiments revealed that the samples consist mainly of the rhombohedral  $\text{Bi}_{2-x}\text{Mn}_x\text{Se}_3$  phase (space group  $R-3m$ ) and some amount of impurity phases  $\text{MnSe}_2$  and  $\text{MnSe}$ , which are proportional to the nominal Mn-content.

The temperature dependences of magnetic susceptibility  $\chi(T)$  of Mn-doped  $\text{Bi}_2\text{Se}_3$  were measured by a SQUID magnetometer in the temperature range of 4.2–300 K at magnetic field of  $H = 0.05$  T. A detectable difference between results of the zero field cooled (ZFC) and field cooled (FC) measurements was not observed. In addition, the anisotropy of magnetic susceptibility appeared to be small. As to the manifestation of the impurity phases  $\text{MnSe}_2$  and  $\text{MnSe}$ , their relative contribution to the magnetic susceptibility was assumed to be negligible because there were no observed evident peculiarities in the measured  $\chi(T)$  at the magnetic transition points intrinsic to these compounds [34].

As seen in figure 1, for temperatures above 20–50 K the experimental  $\chi(T)$  dependence for all the samples obeys a modified Curie–Weiss law



**Figure 1.** Temperature dependence of the reciprocal magnetic susceptibility  $(\chi - \chi_0)^{-1}$  for  $\text{Bi}_{2-x}\text{Mn}_x\text{Se}_3$ . The Curie–Weiss fits are indicated by the solid straight lines.

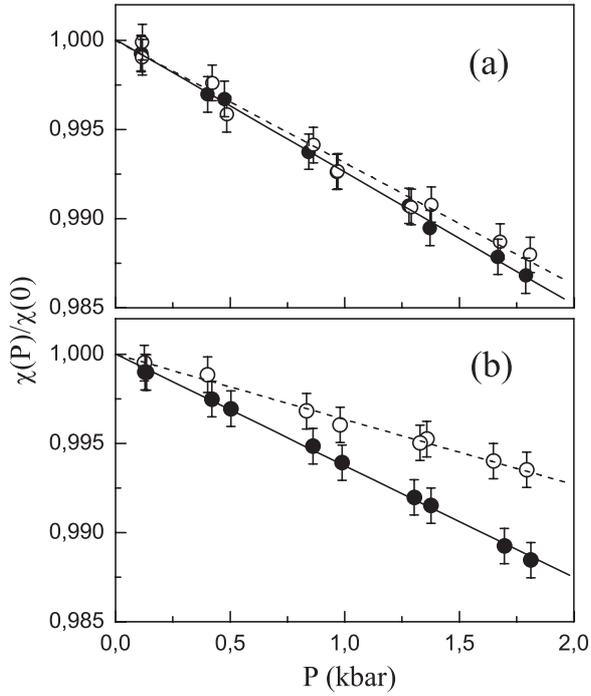
**Table 1.** Curie constant  $C$  (in units  $10^{-4} \text{ K} \cdot \text{emu g}^{-1}$ ), paramagnetic Curie temperature  $\Theta$  (K) and  $\chi_0$  ( $10^{-6} \text{ emu g}^{-1}$ ) in  $\text{Bi}_{2-x}\text{Mn}_x\text{Se}_3$  samples at different nominal Mn content  $x_{\text{nom}}$  and its corrected values  $x_{\text{corr}}$  (see section 3 for details).

$x_{\text{nom}}$	$x_{\text{corr}}$	$C$	$\Theta$	$\chi_0$
0.01	0.01	0.716	$-55.5 \pm 5$	-0.32
0.025	0.017	1.10	$-57.9 \pm 5$	-0.24
0.05	0.05	3.25	$-53.4 \pm 3$	-0.30
0.10 (#1)	0.067	4.76	$-57.8 \pm 3$	-0.30
0.10 (#2)	0.105	7.25	$-54.9 \pm 3$	-0.30
0.10 (#3)	0.140	9.9	$-52.9 \pm 3$	-0.35
0.20	0.20	13.9	$-54.3 \pm 3$	-0.35

$$\chi(T) = \chi_0 + \frac{C}{T - \Theta} \quad (1)$$

where  $\chi_0$  is temperature independent contribution,  $C$  the Curie constant and  $\Theta$  the paramagnetic Curie temperature. The corresponding values of the Curie–Weiss law parameters are collected in table 1.

For the samples of  $\text{Bi}_{2-x}\text{Mn}_x\text{Se}_3$  with nominal Mn content  $x = 0.1$ (#1) and  $x = 0.20$ , the studies of magnetic susceptibility were carried out under helium gas pressure  $P$  up to 2 kbar at fixed temperatures,  $T = 78$  and  $300$  K, using a pendulum-type magnetometer placed directly in the nonmagnetic pressure cell. The magnetometer construction is similar

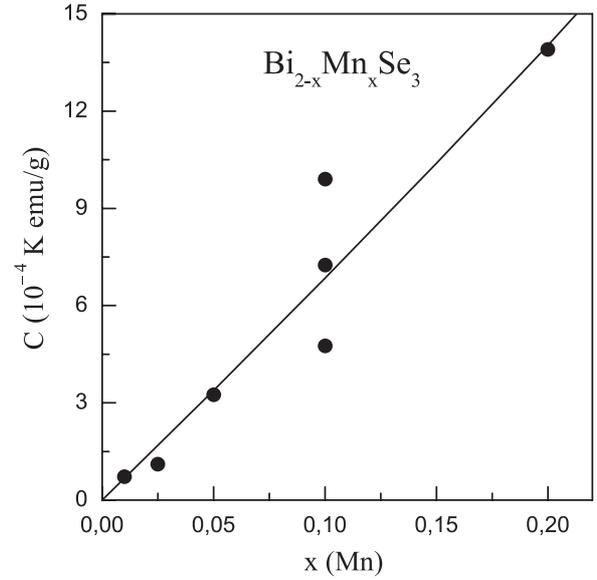


**Figure 2.** Pressure dependencies of the magnetic susceptibility of  $\text{Bi}_{1.9}\text{Mn}_{0.1}\text{Se}_3$  (#1) (a) and  $\text{Bi}_{1.8}\text{Mn}_{0.2}\text{Se}_3$  (b) at  $T = 78$  K (full circles) and 300 K (open circles) normalized to its value at zero pressure.

**Table 2.** Magnetic susceptibility  $\chi$  (in units  $10^{-6}$  emu  $\text{g}^{-1}$ ) and its pressure derivative  $d \ln \chi / dP$  ( $\text{Mbar}^{-1}$ ) for  $\text{Bi}_{2-x}\text{Mn}_x\text{Se}_3$  samples at different temperatures.

Sample	$T$ (K)	$\chi$	$d \ln \chi / dP$
$x = 0.1$ (#1)	78	3.22	$-7.45 \pm 0.5$
	300	1.04	$-6.85 \pm 0.5$
$x = 0.2$	78	10.2	$-6.25 \pm 0.5$
	300	3.58	$-3.5 \pm 0.5$

to that used in [35]. The measured sample with typical sizes of about  $2.5 \times 4 \times 6$  mm<sup>3</sup> was placed inside a small ‘compensating’ coil located at the lower end of the pendulum rod having a length of about 150 mm. The working volume is situated between a cone type pole pieces of an electromagnet in an inhomogeneous magnetic field up to 2 T. The actual diameter of the poles ends and the gap between them were 100 mm and 40 mm, respectively. When the magnetic field is applied, the value of a current through the compensating coil, at which the pendulum comes to its zero position, appears to be a measure of the sample magnetic moment. In order to measure the pressure effects, the pendulum magnetometer is inserted in the cylindrical nonmagnetic pressure chamber with outer and inner diameters of 24 mm and 6 mm, respectively, which is placed in a cryostat. Measurements were performed at a fixed temperature to eliminate the effects on susceptibility of temperature changes upon application or removal of pressure. A detailed description of the device and analysis of the sources of experimental errors will be published elsewhere.



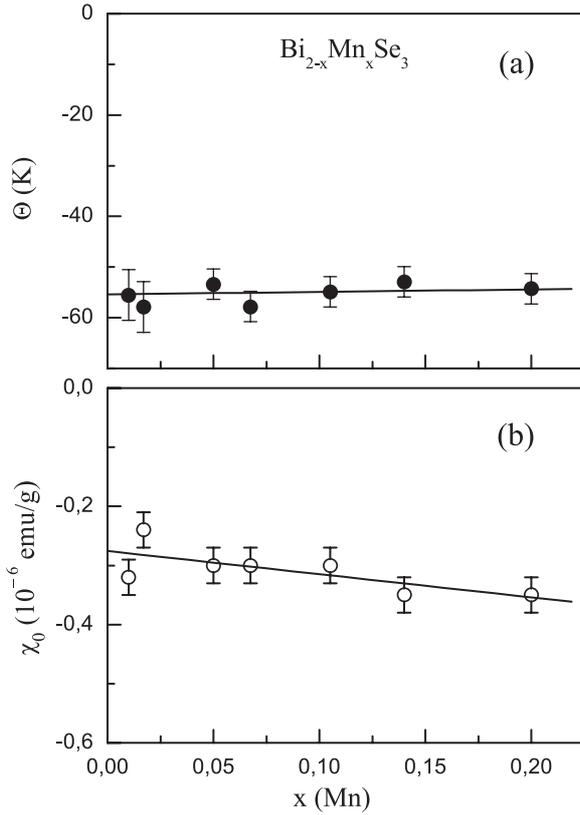
**Figure 3.** Curie constant  $C$  for  $\text{Bi}_{2-x}\text{Mn}_x\text{Se}_3$  samples as a function of Mn content. Solid line corresponds to description of  $C(x)$  assuming effective magnetic moment value of Mn ion to be  $\mu_{\text{eff}} = 5.92 \mu_{\text{B}}$  ( $S = 5/2$ ).

In our case, the relative errors of measurements of  $\chi$  under pressure did not exceed 0.1% for the employed magnetic fields  $H \simeq 1.7$  T. The experimental  $\chi(P)$  dependencies, normalized to the values of  $\chi$  at zero pressure, are presented in figure 2, yielding values of the pressure derivative  $d \ln \chi / dP$  at different temperatures, which are listed in table 2.

### 3. Discussion

As seen in figure 3, the concentration dependence of the Curie constant  $C(x)$  is close to the linear one determined by the value of effective magnetic moment per Mn atom,  $\mu_{\text{eff}} \simeq 5.9 \mu_{\text{B}}$ . This moment corresponds to the spin state of Mn ion  $S = 5/2$ . Some deviation of individual data points (e.g. for  $x = 0.1$ ) are assumed to be due to a difference of the actual content of manganese in the samples,  $x_{\text{corr}}$ , from its nominal value,  $x_{\text{nom}}$ . Resulted from this assumption corrected values of  $x_{\text{corr}}$  are given in table 1 and used to represent concentration dependencies of the Curie–Weiss parameters  $\Theta$  and  $\chi_0$  in figure 4.

As seen in figure 4, the paramagnetic Curie temperature  $\Theta$  in  $\text{Bi}_{2-x}\text{Mn}_x\text{Se}_3$  samples shows negative value, which is weakly dependent on composition and points to a strong antiferromagnetic coupling between magnetic moments of Mn ion. The averaged over the entire range of concentrations value  $\Theta \simeq -55$  K agrees reasonably with the value of  $\Theta \sim -70$  K, which follows from the data of [16] for  $\text{Bi}_{1.97}\text{Mn}_{0.03}\text{Se}_3$  sample, but differs substantially in magnitude from the value  $\Theta \sim -0.5$  K reported in [15] for  $\text{Bi}_{2-x}\text{Mn}_x\text{Se}_3$  compounds with  $x = 0.01$  and 0.02. It should be noted, however, that the latter estimate was based on an analysis of the experimental data on magnetic properties at rather low temperatures,  $2 \leq T \leq 20$  K. At these temperatures the result of [15] could be influenced by possible manifestations of the Kondo and crystal electric field effects in magnetism, as well as the contribution of the impurity ions



**Figure 4.** Values of paramagnetic Curie temperature  $\Theta$  (a) and parameter  $\chi_0$  (b) in  $\text{Bi}_{2-x}\text{Mn}_x\text{Se}_3$  compounds for different Mn content.

of manganese. We believe that this can be a probable explanation for the difference between the estimate of  $\Theta$  in [15] and our results, which were obtained from analysis of  $\chi(T)$  dependences at the higher temperature range (see figure 1(a)). As to the parameter  $\chi_0$ , its values for  $\text{Bi}_{2-x}\text{Mn}_x\text{Se}_3$  compounds correlate with the reported in literature data on magnetic susceptibility of pure  $\text{Bi}_2\text{Se}_3$ ,  $\chi \simeq -0.3 \times 10^{-6} \text{ emu g}^{-1}$  [8] and  $-0.41 \times 10^{-6} \text{ emu g}^{-1}$  [36]. Thus, this parameter can be identified as the intrinsic susceptibility of  $\text{Bi}_2\text{Se}_3$ .

Let us now turn to discuss the effects of pressure on magnetic susceptibility. Neglecting the value of  $\chi_0$  in the Curie–Weiss behaviour (1) of  $\chi(T)$  and assuming the Curie constant  $C$  to be pressure independent, the measured derivative  $d \ln \chi(T)/dP$  is obviously governed by the pressure dependence of the paramagnetic Curie temperature  $\Theta$ :

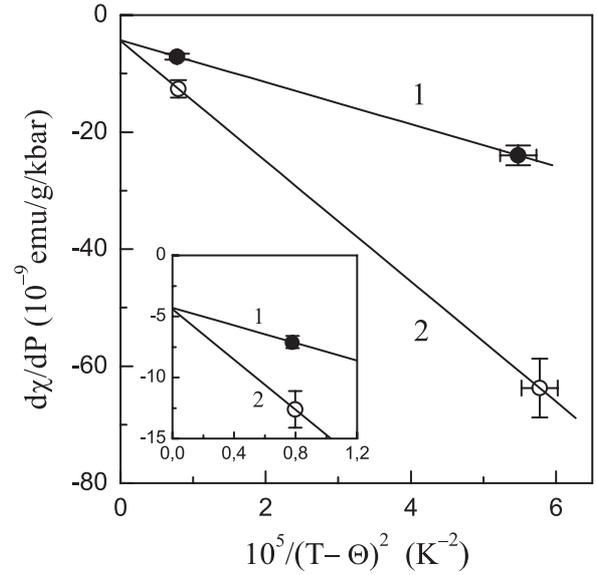
$$\frac{d \ln \chi(T)}{dP} \approx \frac{1}{(T - \Theta)} \frac{d\Theta}{dP}. \quad (2)$$

Then substitution in (2) the experimental data at  $T = 78 \text{ K}$  from table 2 for the sample with  $x = 0.2$  yields the value

$$\frac{d\Theta}{dP} = -0.8 \pm 0.1 \text{ K kbar}^{-1}. \quad (3)$$

In more rigorous analysis the term  $\chi_0$  has to be taken into account and then the pressure effect value  $d\chi/dP \equiv \chi d \ln \chi/dP$  is given by relation

$$\frac{d\chi(T)}{dP} = \frac{d\chi_0}{dP} + \frac{C}{(T - \Theta)^2} \times \frac{d\Theta}{dP}, \quad (4)$$



**Figure 5.** Dependence of the pressure derivative  $d\chi/dP$  for  $\text{Bi}_{1.9}\text{Mn}_{0.1}\text{Se}_3$  #1 (1) and  $\text{Bi}_{1.8}\text{Mn}_{0.2}\text{Se}_3$  (2) versus  $(T - \Theta)^{-2}$ . The inset shows the data near the origin of coordinates on an expanded scale.

being a linear function of  $(T - \Theta)^{-2}$ . The experimental values of  $d\chi/dP$  for the samples with  $x = 0.1$  and  $0.2$  as a function of  $(T - \Theta)^{-2}$  are shown in figure 5. From the slopes of line 1 and 2 in figure 5 one can obtain the values of pressure derivative for paramagnetic Curie temperature to be  $d\Theta/dP = -0.75 \pm 0.05 \text{ K kbar}^{-1}$  and  $-0.74 \pm 0.05 \text{ K kbar}^{-1}$  for  $\text{Bi}_{1.9}\text{Mn}_{0.1}\text{Se}_3$  and  $\text{Bi}_{1.8}\text{Mn}_{0.2}\text{Se}_3$ , respectively. For both compounds, the corresponding pressure derivative of the relative change in  $\Theta$  amounts to

$$d \ln \Theta/dP = 14 \pm 1 \text{ Mbar}^{-1}. \quad (5)$$

The observed substantial pressure effect on  $\Theta$ ,  $d\Theta/dP = -0.75 \text{ K kbar}^{-1}$ , correlates both in sign and magnitude with the reported in [37] value of pressure effect on the Curie temperature  $T_C$  for V-doped  $\text{Sb}_2\text{Te}_3$  compound,  $dT_C/dP = -1.3 \text{ K kbar}^{-1}$ , that suggests a common nature of magnetic interactions in both systems.

Another principal parameter resulted from the data represented in figure 5 is the pressure derivative of the value  $\chi_0$ . It was estimated by extrapolation  $(T - \Theta)^{-2} \rightarrow 0$  to be

$$d\chi_0/dP = -4.3 \pm 0.5 \text{ and } -4.4 \pm 1.5 \quad (6)$$

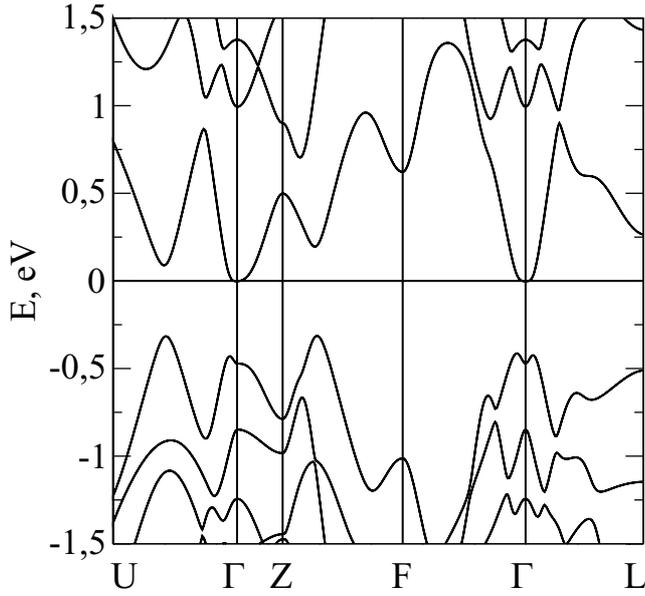
in units of  $10^{-9} (\text{emu g}^{-1}) \text{ kbar}^{-1}$  for the samples with  $x = 0.1$  and  $0.2$ , respectively. Using this estimate and the value  $\chi_0 \simeq -0.3 \times 10^{-6} \text{ emu g}^{-1}$  from table 1, one obtains

$$d \ln \chi_0/dP = 14 \pm 3 \text{ Mbar}^{-1}. \quad (7)$$

As can be seen from equations (5) and (7), both  $\Theta$  and  $\chi_0$  parameters show significant pressure effects, which are similar in magnitude and correlate with the experimentally observed changes under pressure of the carrier density  $n$

$$d \ln n/dP = 15 \div 20 \text{ Mbar}^{-1}, \quad (8)$$

resulted from the initial slope of the  $n(P)$  dependence in  $\text{Bi}_2\text{Se}_3$  [32, 33].



**Figure 6.** Band structure of  $\text{Bi}_2\text{Se}_3$  along the symmetry directions of the rhombohedral Brillouin zone. The Fermi level is indicated by the horizontal line at  $E = 0$ .

It is quite curious that the observed Curie temperature  $\Theta$  itself and its pressure derivative do not notably depend on the Mn concentration. This allows to presume that Mn doping effects have the single-site impurity origin, e.g. of the Kondo type with  $T_K \approx -\Theta/4 \sim 10\text{--}15$  K. The Kondo temperature  $T_K$  is given by [38, 39]:

$$k_B T_K \approx D \exp[-1/JN(E_F)], \quad (9)$$

where  $D$  is the effective band width,  $J$  the interaction parameter between Mn-moment and conduction electron spins,  $N(E_F)$  density of electronic states at the Fermi energy  $E_F$ . Then, keeping in mind the strong pressure dependence of the carrier density (8), one can assume the dominant contribution in the pressure effect on  $T_K(\Theta)$  to be due to the increasing under pressure the density of states  $N(E_F)$ , which is proportional to  $n^{1/3}$  for the parabolic band.

An alternative mechanism for the observed effects of manganese doping involves the formation of the Mn-enriched clusters with similar magnetic properties instead of a uniform distribution of the manganese atoms throughout the crystal. For this case, the paramagnetic Curie temperature  $\Theta$  in the cluster can be determined by an indirect exchange interaction between Mn-moments mediated by conduction electrons. In the framework of the conventional RKKY approach [40], the main functional dependence of  $\Theta$  may be given by:

$$\Theta \propto J^2 n^{4/3} F(n) \quad (10)$$

where  $J$  is the exchange coupling parameter between Mn-moments and conduction electron spins,  $F(n)$  the RKKY function. As seen from equation (10), the strong pressure dependence of  $\Theta$  could be governed by the pressure effect on the carrier density  $n(P)$ .

It should be admitted that in order to make a choice between the above proposed approaches, we need a more rigorous

quantitative analysis, which represent a challenging task and is not the subject of this paper.

Let us conclude by discussing the derived pressure effect on the parameter  $\chi_0$ , which can be identified as the intrinsic susceptibility of  $\text{Bi}_2\text{Se}_3$ ,  $\chi$ . It can generally be represented as the sum:

$$\chi = \chi_P + \chi_{VV} + \chi_{LP} + \chi_{ion}. \quad (11)$$

Here  $\chi_P$  is the Pauli spin susceptibility,  $\chi_{VV}$  a generalization of the Van Vleck orbital paramagnetism,  $\chi_{LP}$  the Landau–Peierls diamagnetism and  $\chi_{ion}$  the Langevin diamagnetism of closed ion shells. For a parabolic band, the Pauli paramagnetism and Landau–Peierls diamagnetism are described by well known relations [41]

$$\chi_P = \frac{4m^* \mu_B^2}{h^2} (3\pi^2 n)^{1/3}, \quad \chi_{LP} = -\frac{4m_0^2 \mu_B^2}{3m^* h^2} (3\pi^2 n)^{1/3}, \quad (12)$$

where  $m^*$  is the effective mass of band electrons. Further, using for  $\text{Bi}_2\text{Se}_3$  the estimate of  $m^* = 0.11 \div 0.18 m_0$  [42–44], one can conclude that  $\chi_{LP} \gg \chi_P$ . The value of orbital susceptibility was calculated to be  $\chi_{VV} \approx 0.145 \times 10^{-6} \text{ emu g}^{-1}$  [27], and the Langevin term,  $\chi_{ion} \approx -0.11 \times 10^{-6} \text{ emu g}^{-1}$ , results from the data for the ion susceptibilities of  $\text{Bi}^{5+}$  and  $\text{Se}^{6+}$  [45]. Based on the above estimates and the measured value of magnetic susceptibility for  $\text{Bi}_2\text{Se}_3$ ,  $\chi \approx -0.30 \times 10^{-6} \text{ emu g}^{-1}$ , we have evaluated the Landau–Peierles term to be

$$\chi_{LP} \approx -0.34 \times 10^{-6} \text{ emu g}^{-1}. \quad (13)$$

Because the orbital contribution is weakly dependent on pressure and the Pauli paramagnetism is relatively small, it is reasonable to assume that the Landau–Peierles term gives the main contribution to the measured pressure effect, i.e.  $d\chi_0 \approx d\chi_{LP}$  and

$$\frac{d \ln \chi_{LP}}{dP} \equiv \frac{1}{\chi_{LP}} \frac{d\chi_0}{dP} \approx 12.5 \text{ Mbar}^{-1}. \quad (14)$$

As it follows from equation (12),

$$\frac{d \ln \chi_{LP}}{dP} = \frac{1}{3} \frac{d \ln n}{dP} - \frac{d \ln m^*}{dP}. \quad (15)$$

In order to estimate the pressure dependence of  $m^*$ , we have carried out the calculations of the volume dependent band structure of  $\text{Bi}_2\text{Se}_3$ , which is presented in figure 6.

The calculations were performed by employing a relativistic full-potential LMTO method (FP-LMTO, RSPt code [46–48]) within the generalized gradient approximation [49] of the density functional theory. The effective mass of electrons was obtained using a parabolic band model at the conduction band minimum by a parabolic fitting to the band dispersion  $E(k)$  along different directions around the  $\Gamma$  symmetry point of the Brillouin zone (see figure 6) and employing the relation  $m^* = 1/(\partial^2 E/\partial k^2)$  (in atomic units). The calculated average value of the effective mass,  $m^* \approx 0.12 m_0$ , is found to increase with volume  $V$ , and the corresponding volume derivative of  $m^*$  appeared to be  $d \ln m^*/d \ln V \approx 3$ . The obtained values are assumed to be valid for small carrier density inherent to real samples of the  $\text{Bi}_2\text{Se}_3$  compound.

Taking into account the bulk modulus value  $B \simeq 0.5$  Mbar for  $\text{Bi}_2\text{Se}_3$  [50], we have estimated the pressure derivative of  $m^*$  in (15) as  $d \ln m^*/dP \simeq -6$  Mbar $^{-1}$ . By substituting this estimate to equation (15), together with the first term,  $(1/3)d \ln n/dP = 5 \div 7$  Mbar $^{-1}$  [32, 33], we have obtained a reasonable agreement of the model (15) with the experimental result (14).

#### 4. Summary

In conclusion, the temperature dependence of magnetic susceptibility for  $\text{Bi}_{2-x}\text{Mn}_x\text{Se}_3$  compounds was measured for the concentration range up to  $x = 0.2$ . For all the samples, the experimental data obey a Curie–Weiss law with the effective magnetic moment of Mn ions corresponding to the spin value  $S = 5/2$  and antiferromagnetic coupling of the moments. For Mn-rich samples, the dependence of magnetic susceptibility on the pressure was measured for the first time.

The most surprising result is that the observed Curie–Weiss parameters,  $\Theta$  and  $\chi_0$ , and their available pressure derivatives do not depend notably on the Mn concentration. This behaviour suggests that the effects of Mn doping are the single-site impurity origin of the Kondo type or these effects involve the formation of the Mn-enriched clusters of the similar magnetic properties instead of a uniform distribution of the manganese atoms throughout the crystal. In both approaches, the large pressure effects in magnetism of the system points to their strong correlation with the literature data on the pressure dependence of the carrier density in  $\text{Bi}_2\text{Se}_3$ .

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